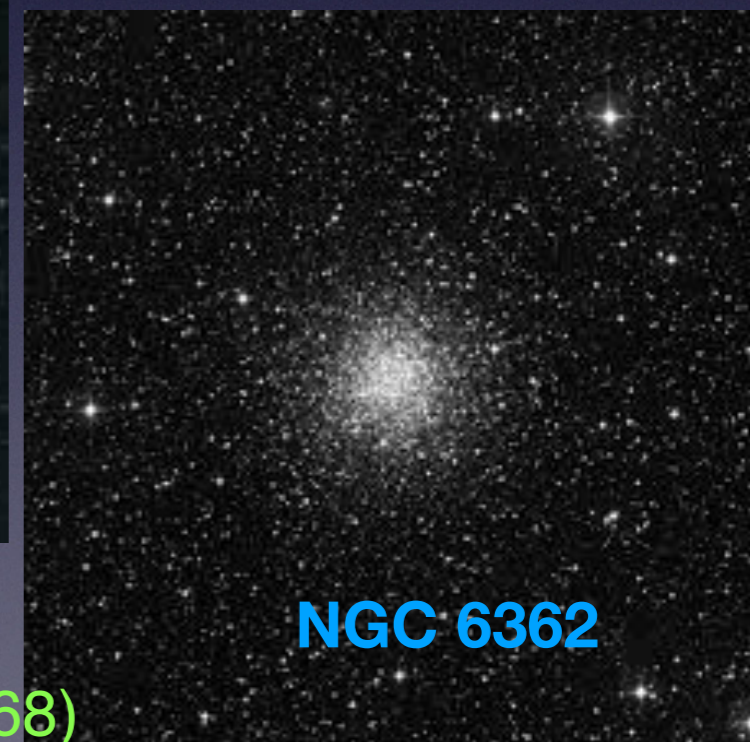




## *Evaluating the Merger Rate of Binary Black Holes Using the Milky Way Globular Clusters*



Kritos & Cholis, PRD 102, 8, 083016 (2020) (arXiv:2007.02968)

**Midwest Relativity Meeting 2020**

**Ilias Cholis, 10/24/2020**

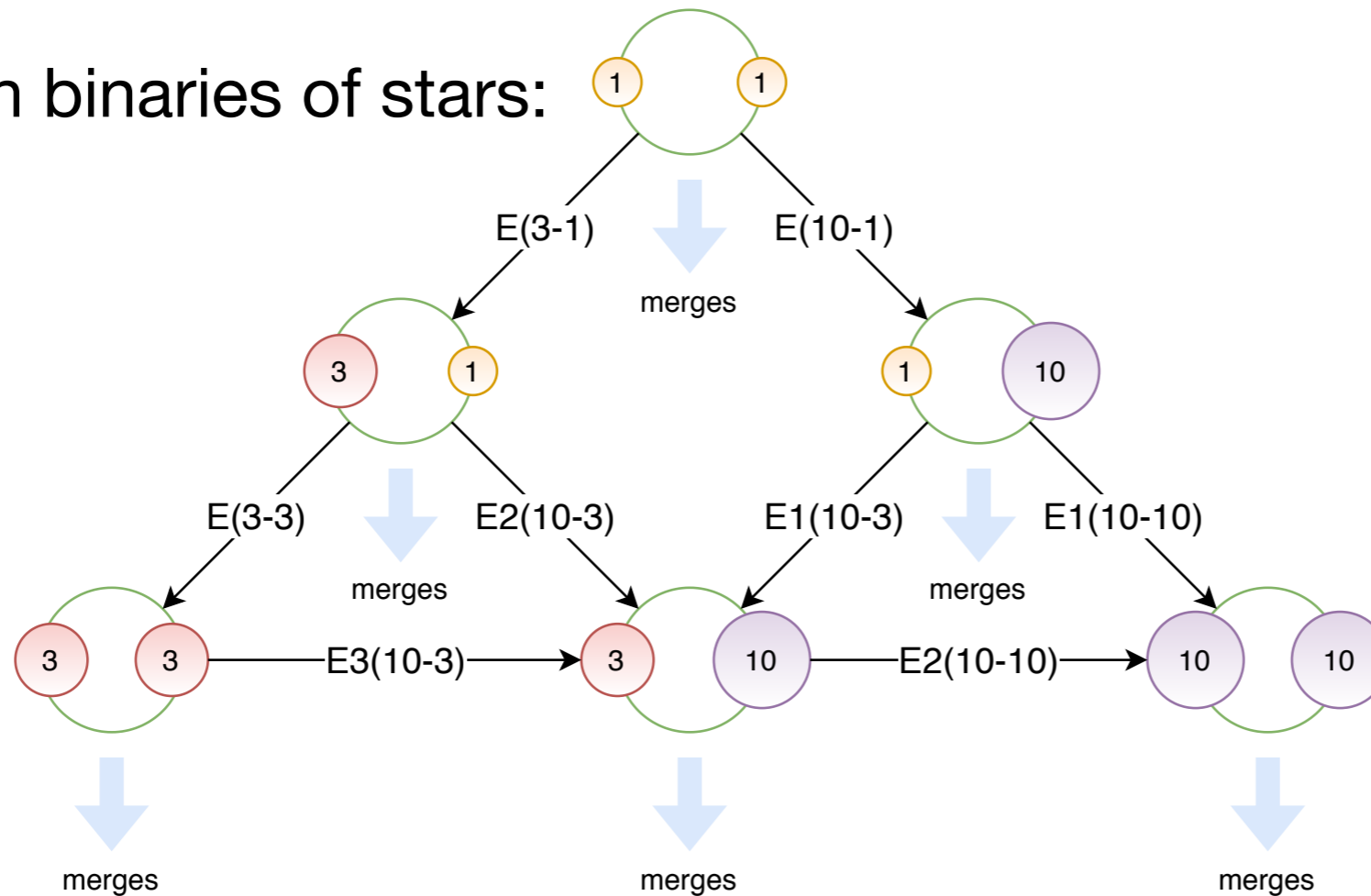
## ***The Centers of Globular Clusters***

Very dense stellar environments where multiple dynamical interactions take place. Most stars are in binaries.

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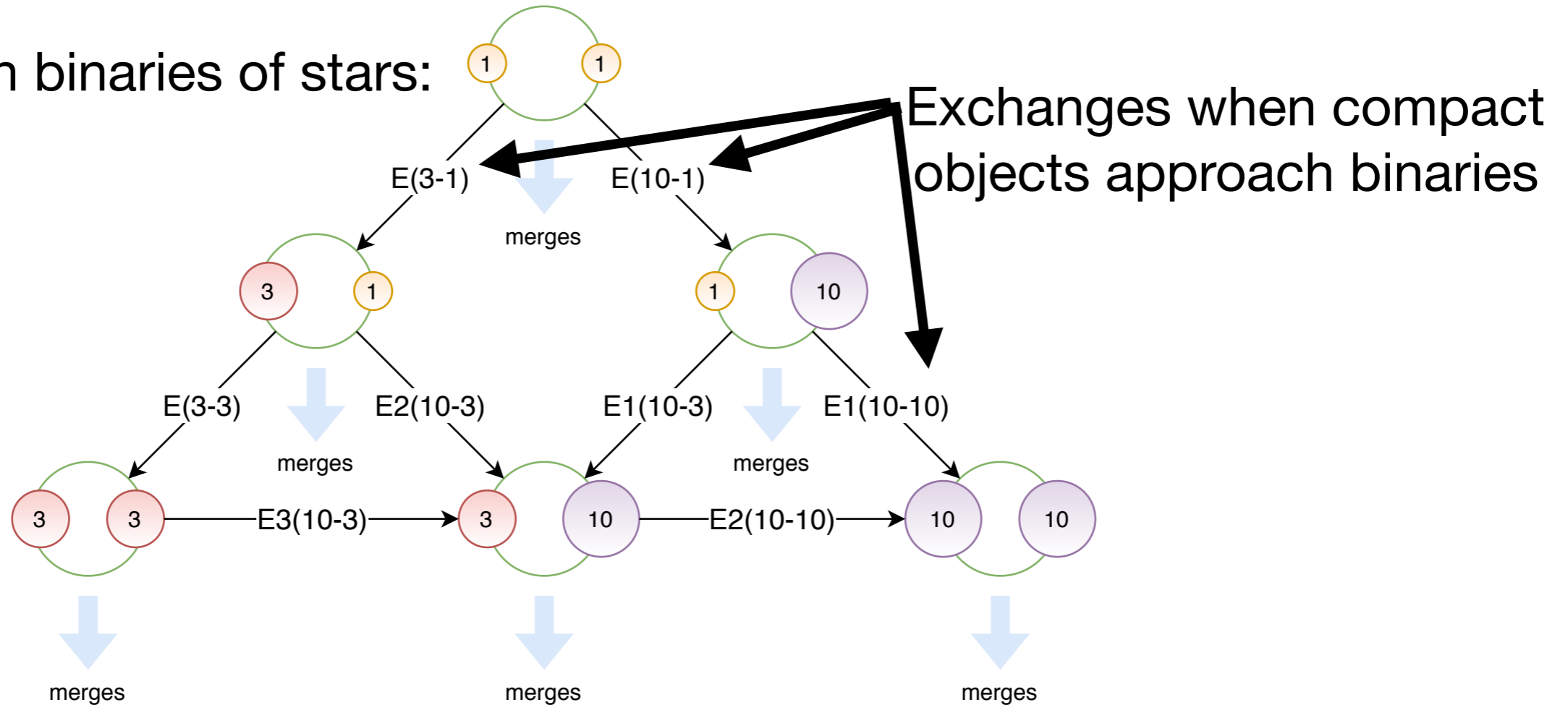
Start with binaries of stars:



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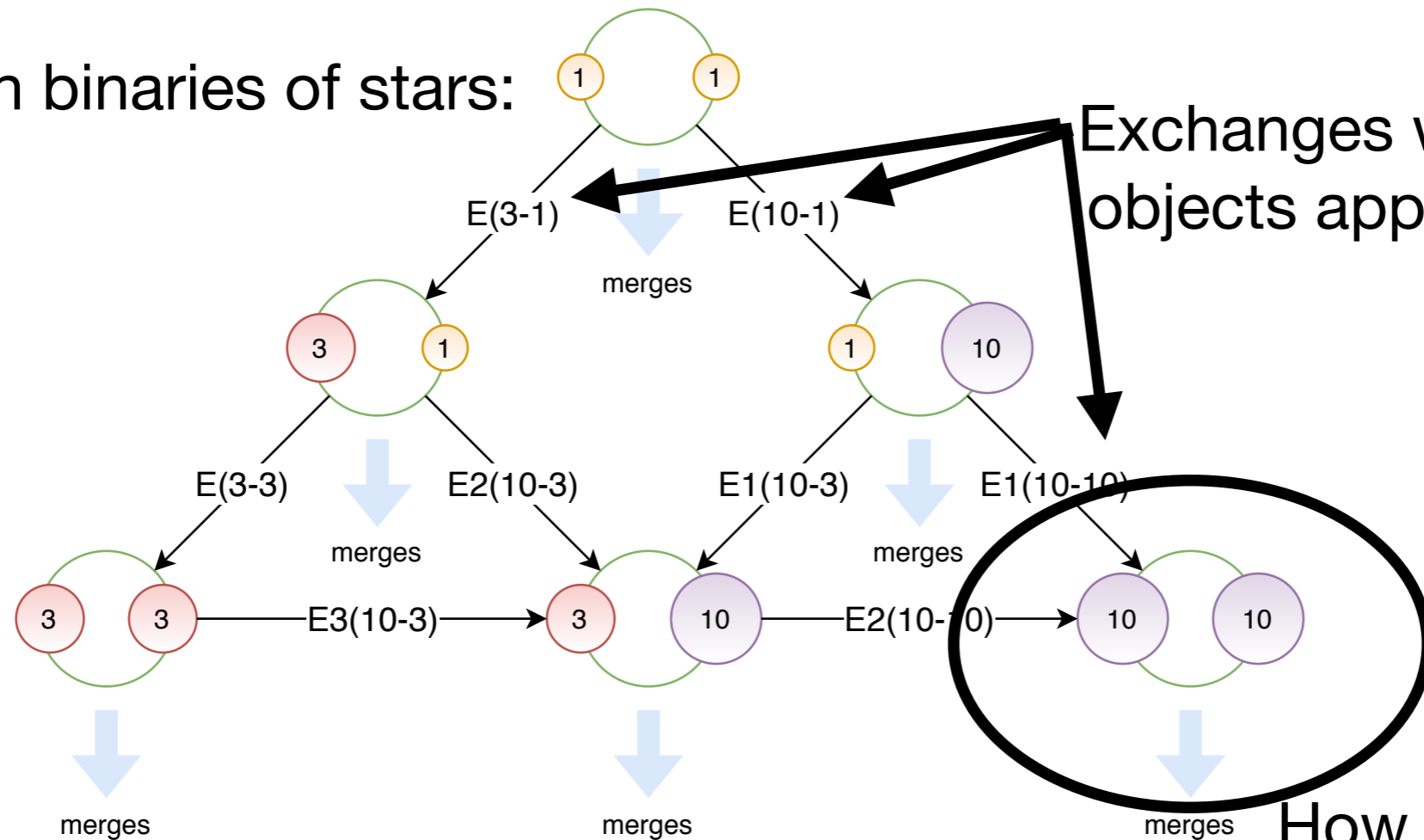
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# The Centers of Globular Clusters

Very dense stellar environments where multiple dynamical interactions take place. Most stars are in binaries.

Start with binaries of stars:



Exchanges when compact objects approach binaries

Focus on the BH-BH binaries.

How the binaries evolve while interacting with third-bodies leading to their merger

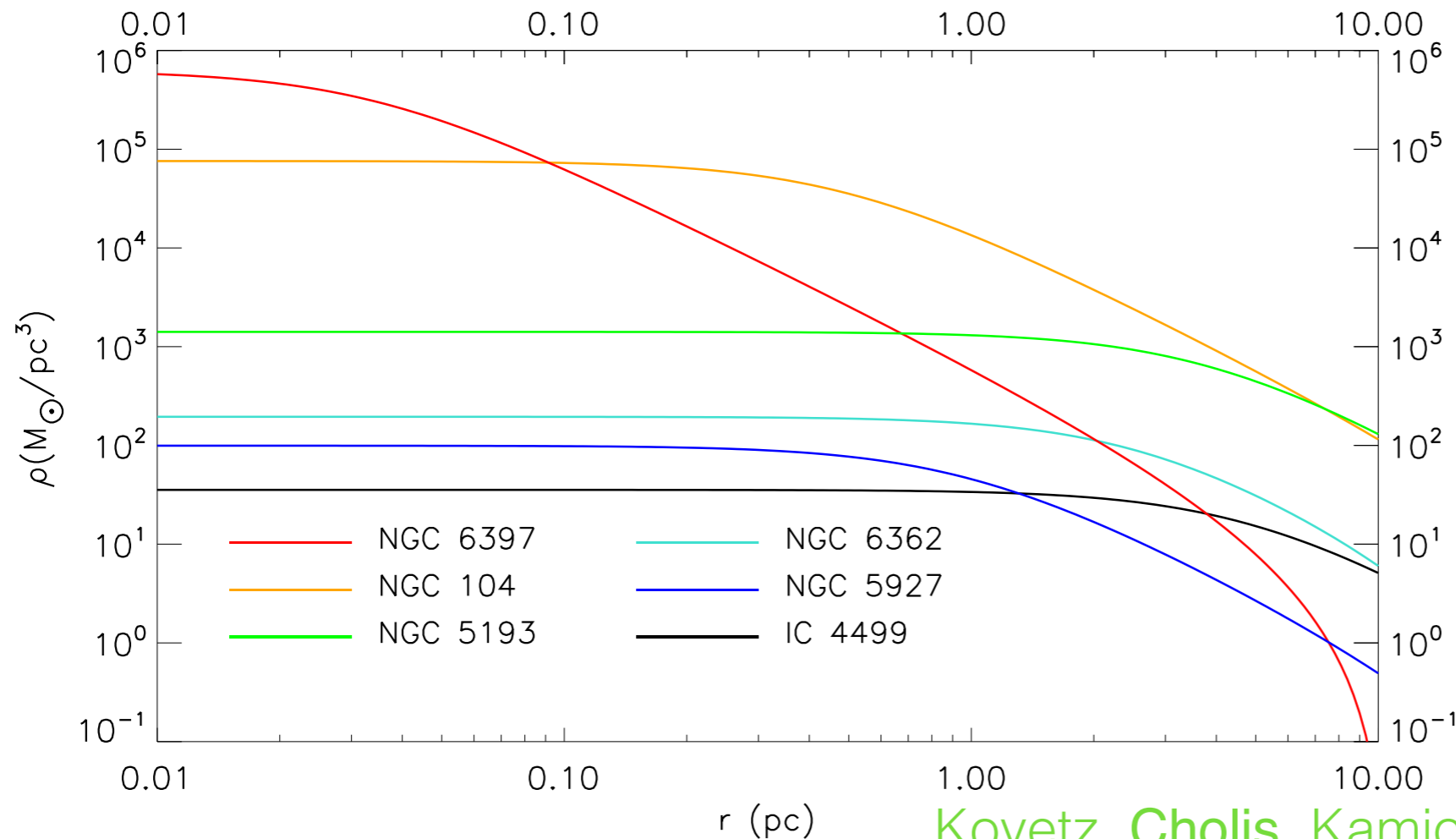
Another channel to produce merging BH-BH binaries comes from direct captures.

# The Centers of Globular Clusters

Very dense stellar environments where multiple dynamical interactions take place. Most stars are in binaries.

*There are about 150 Globular Clusters in the Milky Way. Significant enough sample for us to evaluate the BH-BH merger rates from globular clusters as a population.*

*Most of which have well measured properties:*



Data from Harris  
GC catalogue

Kovetz, Cholis, Kamionkowski, Silk PRD 2018

Assume for simplicity that all BHs are  $10 M_{\odot}$

We can estimate the mass fraction of the GC that ends up in BHs,

$$f_{BH} \simeq \frac{1}{3} \frac{1}{M_{GC}} \int_{25M_{\odot}}^{120M_{\odot}} d\mathcal{M} \mathcal{M} \xi(\mathcal{M}) \approx 0.03$$

Kroupa Initial Mass Function mass of stars

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From that we can estimate the maximum number of BHs that were born at any point in a GC,

$$N_{BH}^{\max} = f_{BH} \frac{M_{GC}}{10 M_{\odot}}$$

and including the natal kicks that are responsible for roughly 90% of the BHs escaping the clusters, the maximum number of retained BHs,

$$N_{BH}^{\text{ret-max}} = f_{\text{ret}} \times N_{BH}^{\max}$$



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**BHs that could end up in binaries inside the GCs**

	GC	$\frac{r_c}{1\text{pc}}$	$c$	$\frac{r_{pl}}{1\text{pc}}$	$\log_{10} \left( \frac{\rho_0}{1M_{\odot}/\text{pc}^3} \right)$	$\frac{M_{GC}}{10^5 M_{\odot}}$	$N_{BH}^{\text{ret-max}}$	
From that we can	47 Tuc	0.47	2.07	0.73	4.88	38.2	1145	at
any point in a GC	$\omega$ Cen	3.60	1.31	5.57	3.15	49.3	1477	
	M15	0.42	2.29	0.66	5.05	68.7	2060	
	M22	1.24	1.38	1.92	3.63	7.30	219	
and including the	NGC 6362	2.50	1.09	3.88	2.29	1.29	38	
escaping the clus	NGC 5946	0.25	2.50	0.38	4.68	9.44	283	
	M 30	0.14	2.50	0.22	5.01	3.80	113	
	Terzan 5	0.32	1.62	0.50	5.14	7.51	225	
$N_{BH}^{\text{ret-max}} = f_{\text{ret}} \times$	Pal 2	1.35	1.53	2.09	4.06	36.8	1103	
	NGC 6139	0.44	1.86	0.69	4.67	11.7	351	
	NGC 2808	0.70	1.56	1.08	4.66	22.1	661	
	NGC 5286	0.95	1.41	1.48	4.10	10.6	319	
	NGC 6316	0.51	1.65	0.80	4.23	4.08	122	

Relying on work by Banerjee et al. MNRAS 2010 and O'Leary et al. ApJ 2006 we take that only about 1-10% of these BHs will make hard binaries (binaries that will not break from interactions with regular stars). Those are binaries that should have a semi-major axis of at most,

$$a_h = \frac{G m_{BH}}{4 \sigma^2} \simeq 5.58 \times \left( \frac{m_{BH}}{10 M_{\odot}} \right) \left( \frac{20 \text{ km/s}}{\sigma_{\text{star}}} \right)^2 \text{ AU}$$

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After every encounter the hard BH-BH binaries will typically loose :

$$\frac{\langle \Delta E_b \rangle}{E_b} \simeq 0.12 \times \left( \frac{H}{15} \right) \left( \frac{m_{\text{star}}}{1 M_\odot} \right) \left( \frac{10 M_\odot}{m_{BH}} \right)$$

Hardening rate, it takes ~15 encounters for the BH-BH to reach the point where the GW emission will dominate its evolution (see Sesana et al ApJ 2006).

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Hardening rate, it takes ~15 encounters for the BH-BH to reach the point where the GW emission will dominate its evolution (see Sesana et al ApJ 2006). The combined effects of third-body encounters and GW emission on the semi-major axis and the eccentricity of the binaries is given by:

$$\dot{a} = - \frac{G H \rho_{\text{star}}}{\sigma_{\text{star}}} a^2 - \frac{128}{5} \frac{G^3 m_{BH}^3}{c^5 a^3} F(e)$$

$$\dot{e} = + \frac{G H K \rho_{\text{star}}}{\sigma_{\text{star}}} a - \frac{608}{15} \frac{G^3 m_{BH}^3}{c^5 a^4} D(e)$$

Peters GW emission terms

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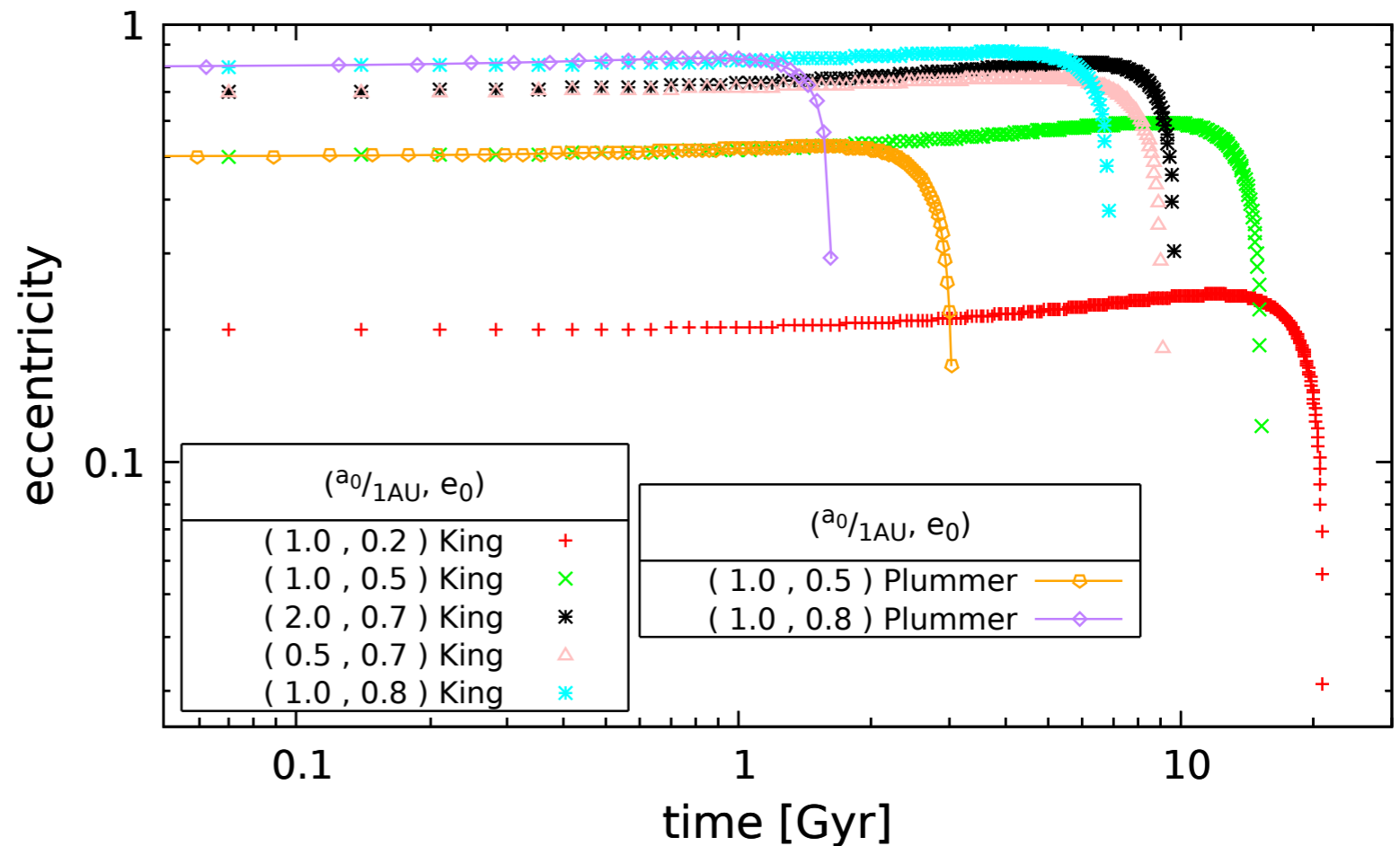
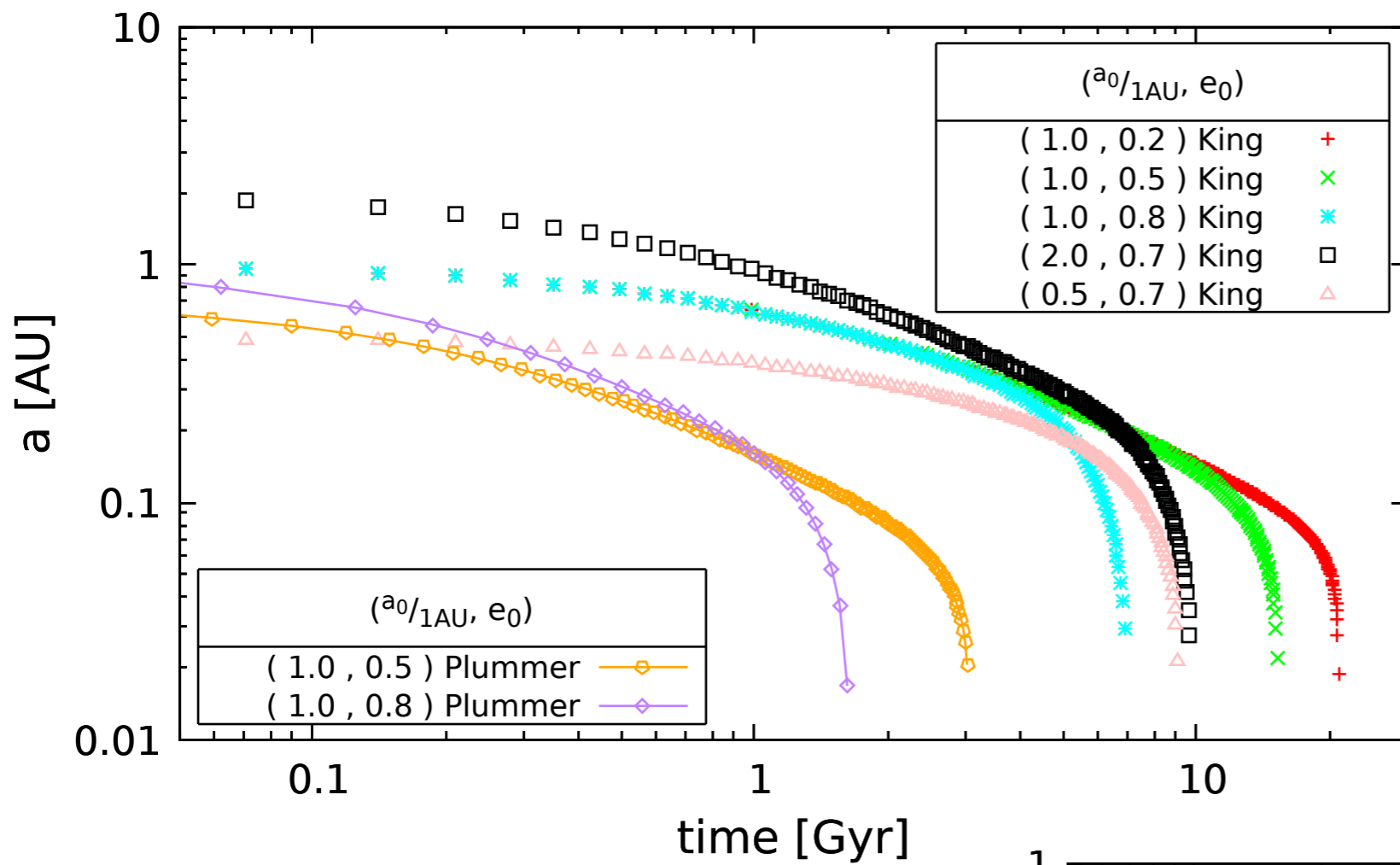
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**Environment-dependent terms**

Peters GW emission terms

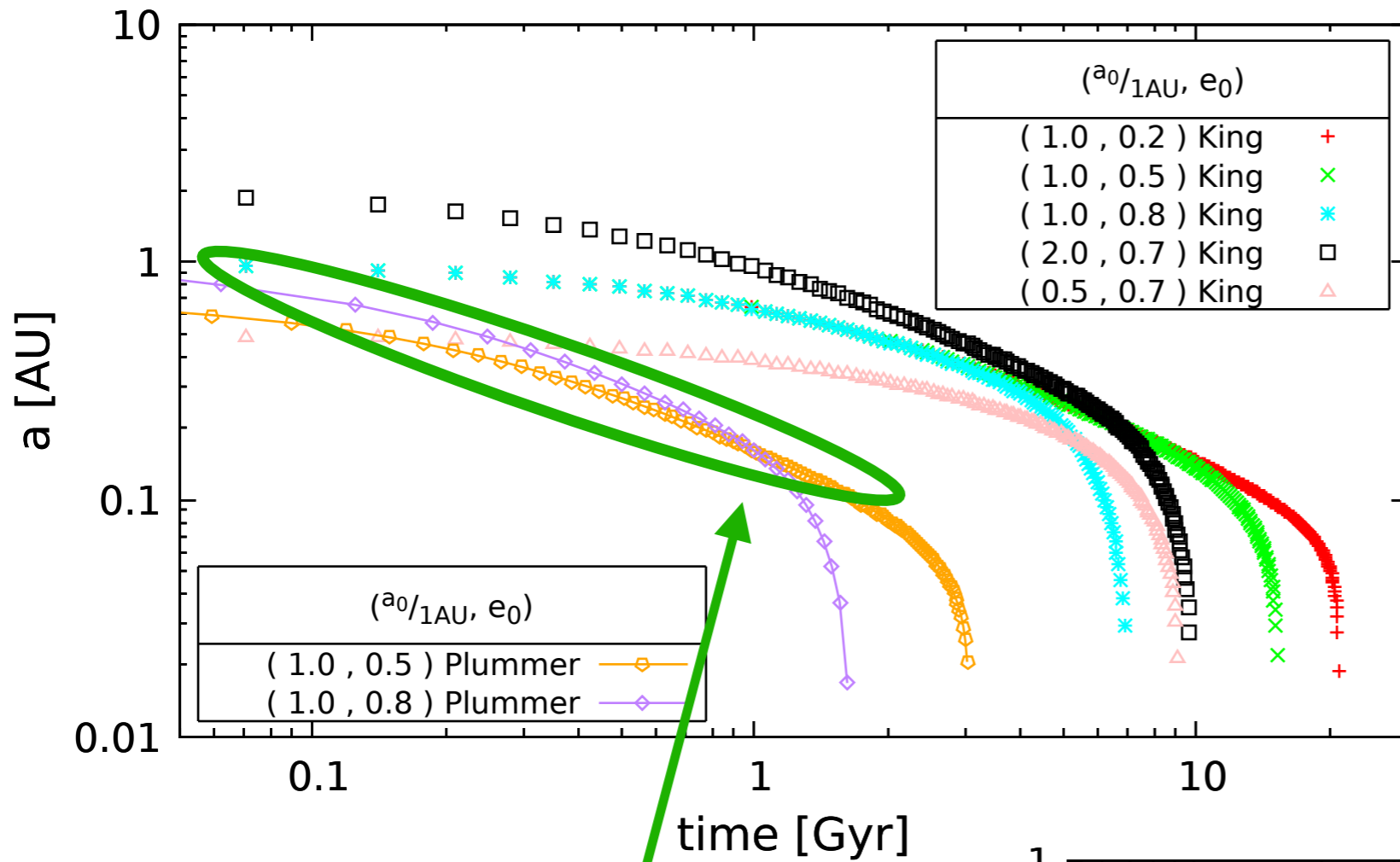
## Example: Environment of 47 Tuc (NGC 104)

The evolution of a 10 + 10 solar masses BH-BH binary, for different initial assumptions on its semi-major axis and eccentricity

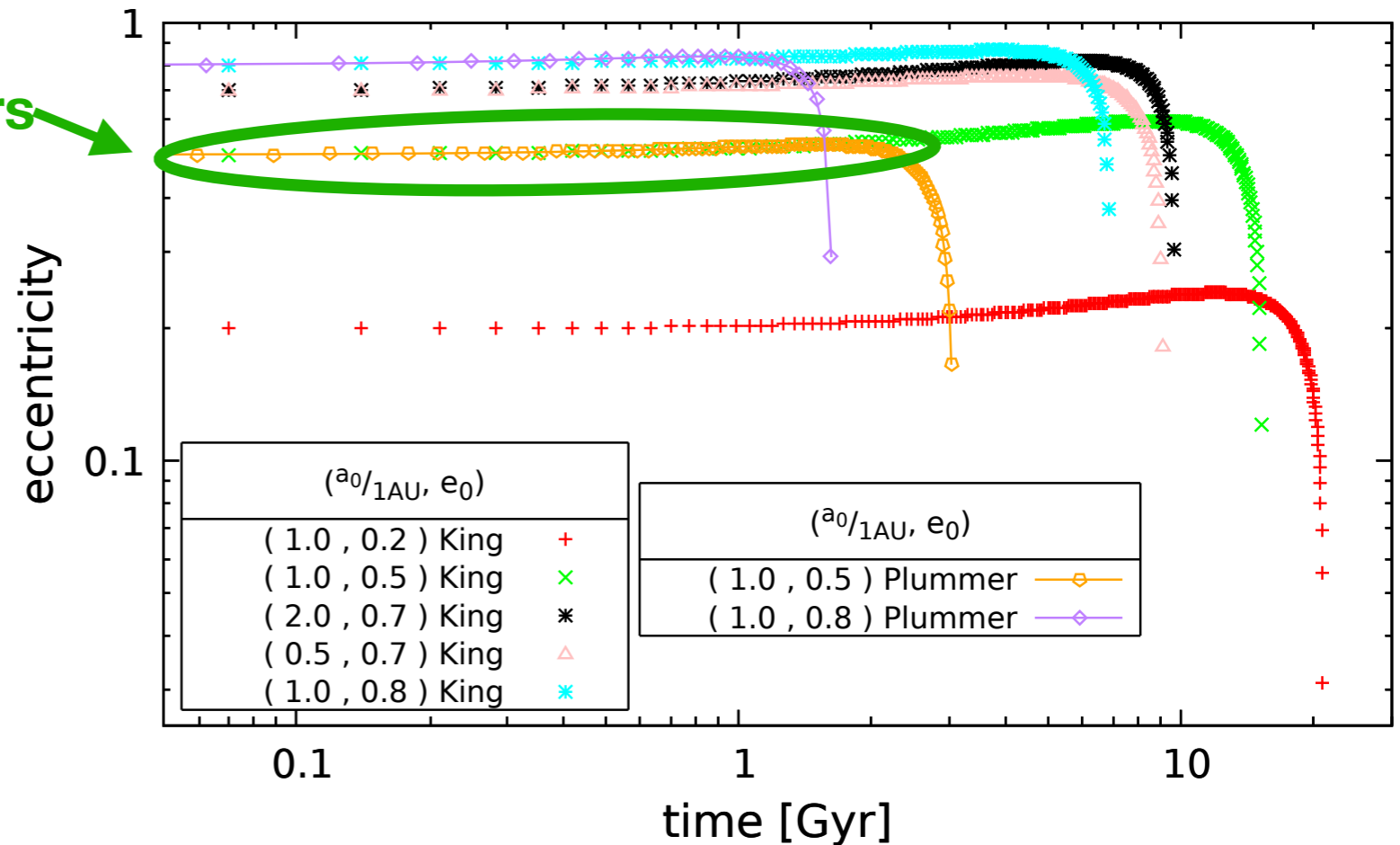


## Example: Environment of 47 Tuc (NGC 104)

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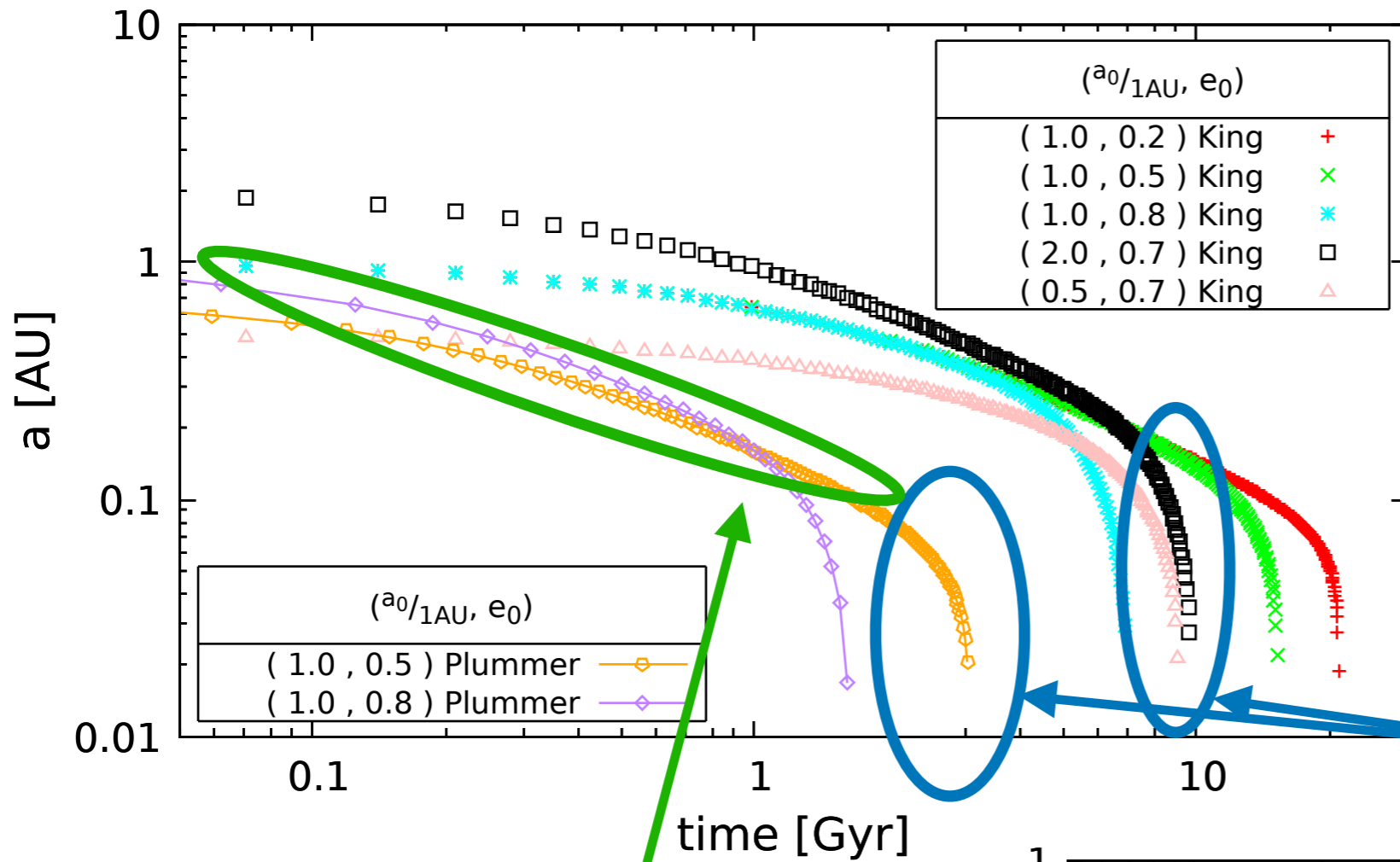


third-body interactions with stars



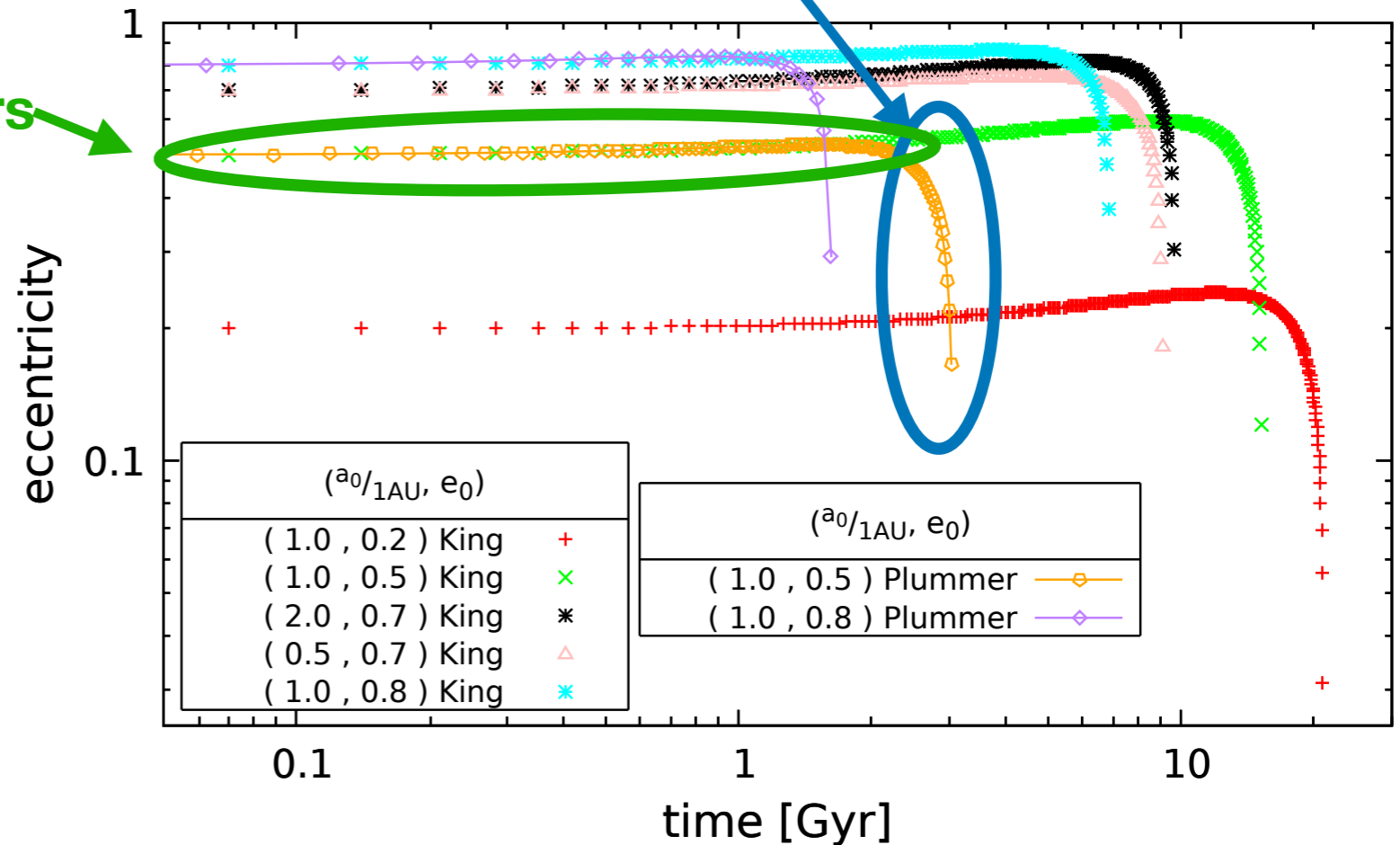
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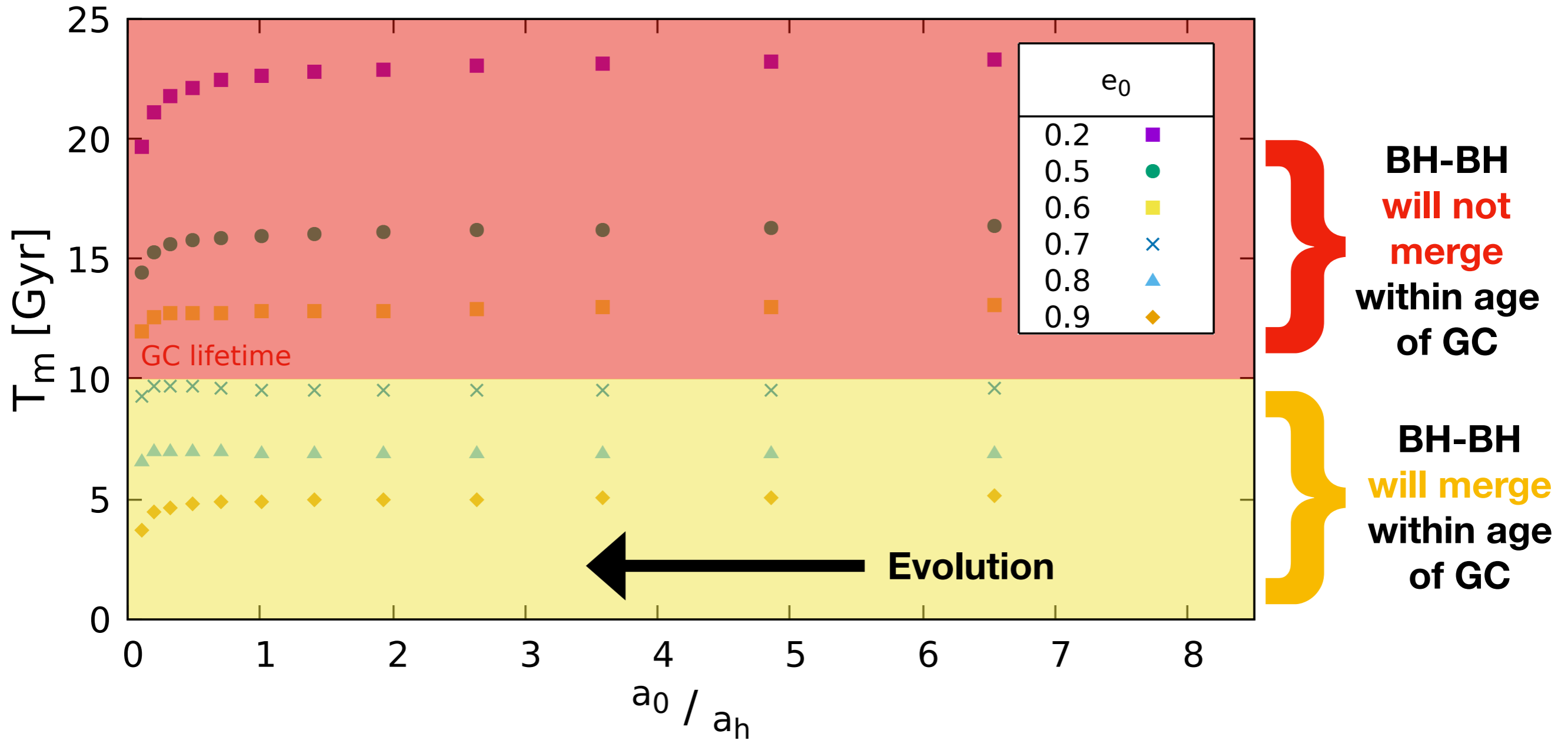
BH-BH binaries with **high initial eccentricities** will merge faster and within the age of the GC. Instead the initial semi-major axis assumptions are of little significance.



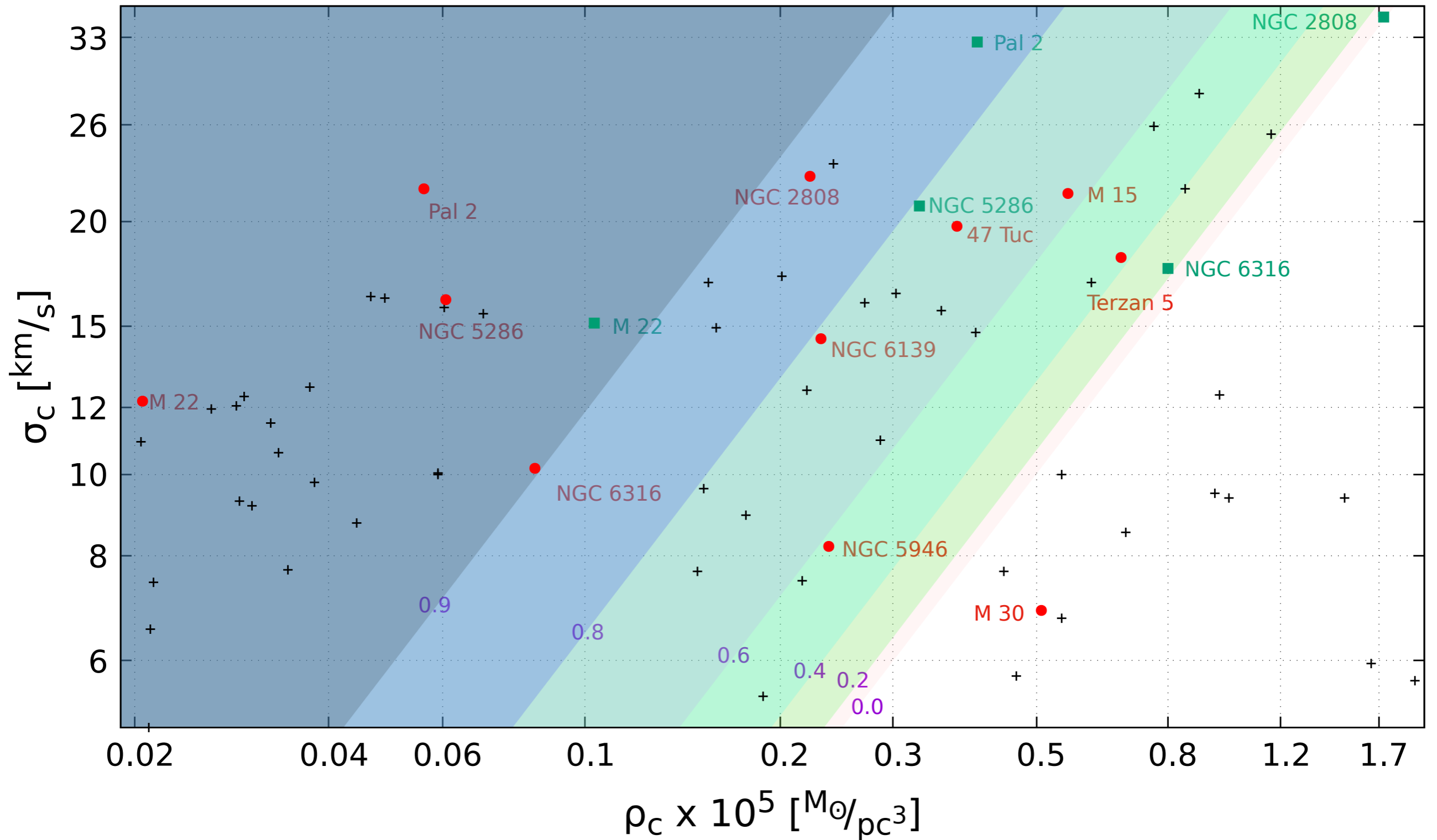
GW emission dominance



# Example: Environment of 47 Tuc (NGC 104)

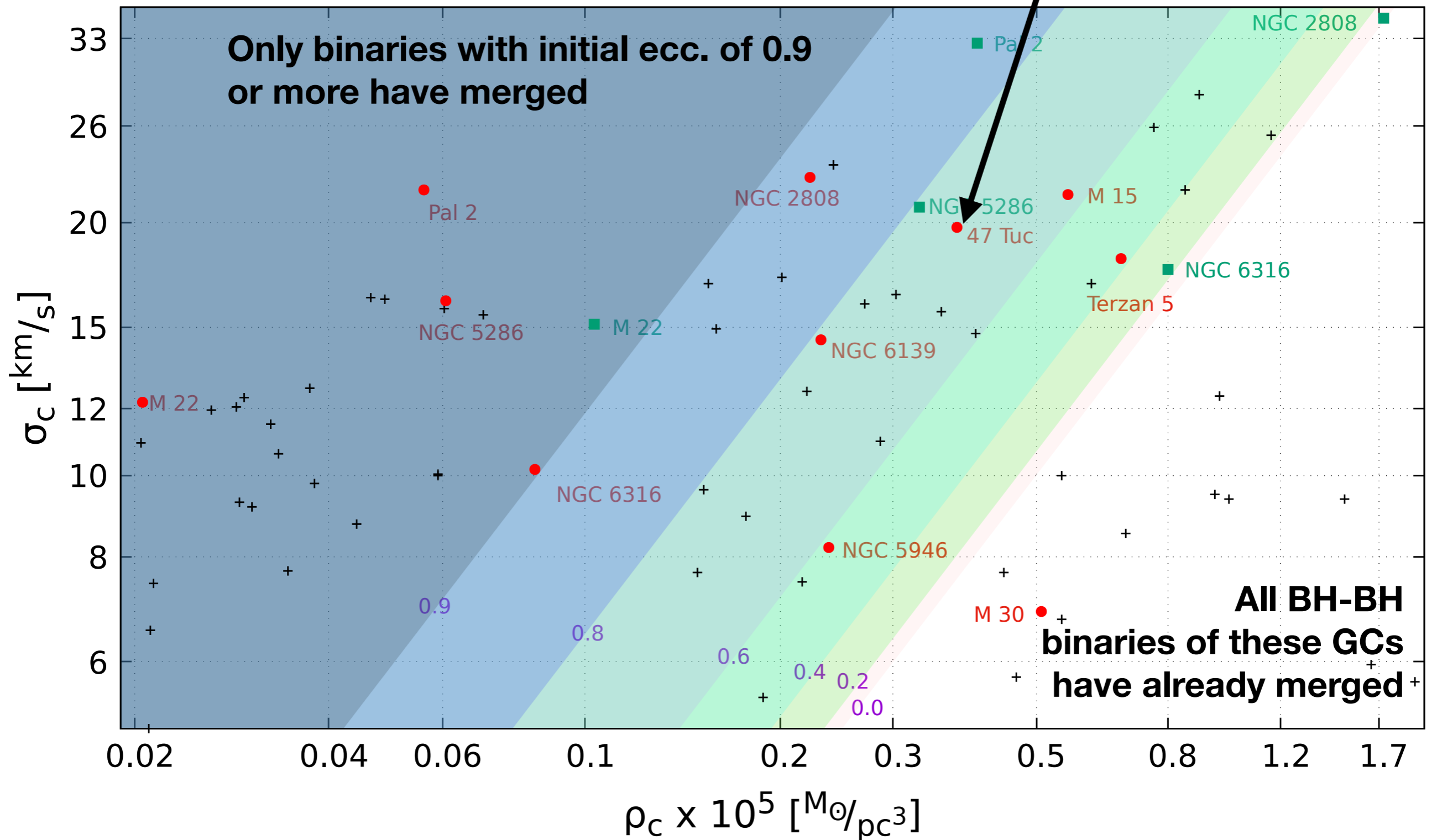


# Connecting to the observed properties of Milky Way Globular Clusters



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About half of the BH-BH binaries (ecc. > 0.7) inside 47 Tuc have already merged



## Concluding

- Once averaging over the Milky Way sample third-body soft interactions give an averaged rate of  $2 - 4 \times 10^{-10} \text{yr}^{-1}$  per cluster.
- Accounting for direct capture events adds  $0.3 - 5 \times 10^{-11} \text{yr}^{-1}$  per cluster.
- The BBH mergers inside globular clusters can be responsible for  $\sim 100$  mergers per year up to redshift of  $z=1$ .
- We have made a connection between the observed properties of clusters and the expected BH-BH merger events.
- Further observations of cluster properties on their total mass, density and velocity profiles and a better modeling of their cosmological distribution will allow us to more accurately evaluate those environments' contribution to the coalescence events observable from GW detectors.